

EXPRESS LETTER

Continuously Doping $\text{BiM}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ into Electron-Doped Superconductor by CaH_2 Annealing Method

To cite this article: Jin Zhao et al 2022 Chinese Phys. Lett. 39 077403

View the article online for updates and enhancements.

You may also like

- Critical temperature and critical current density of hydrogen-doped SmFeAsO epitaxial films fabricated by thermal annealing with binary hydrides Hidenori Hiramatsu, Jumpei Matsumoto and Hideo Hosono
- Evolution of incommensurate superstructure and electronic structure with Pb substitution in (Bi₂, Pb_)Sr₂CaCu₂O₈₊ superconductors Jing Liu, , Lin Zhao et al.
- <u>The synthesis of magnetically exchange</u> <u>coupled CoFe₂O₄/CoFe₂ composites</u> <u>through low temperature CaH₂ reduction</u> Weimin Li, Yashu Liu and Yuhang Guo

Continuously Doping $Bi_2Sr_2CaCu_2O_{8+\delta}$ into Electron-Doped Superconductor by CaH_2 Annealing Method

Jin Zhao(赵金)^{1,2†}, Yu-Lin Gan(甘渝林)^{1,2†}, Guang Yang(杨光)^{1,2†}, Yi-Gui Zhong(钟益桂)^{1,2,3}, Cen-Yao Tang(唐岑瑶)^{1,2}, Fa-Zhi Yang(杨发枝)^{1,2}, Giao Ngoc Phan¹, Qiang-Tao Sui(随强涛)^{1,2}, Zhong Liu(刘忠)¹, Gang Li(李岗)¹, Xiang-Gang Qiu(邱祥冈)¹, Qing-Hua Zhang(张庆华)¹, Jie Shen(沈洁)¹, Tian Qian(钱天)^{1,4}, Li Lu(吕力)¹, Lei Yan(颜雷)¹, Gen-Da Gu(顾根大)⁵, and Hong Ding(丁洪)^{1,2,4*}

¹Beijing National Laboratory for Condensed Matter Physics, and Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China

²School of Physics, University of Chinese Academy of Sciences, Beijing 100190, China

³Institute for Solid State Physics, University of Tokyo, Chiba 277-8581, Japan

⁴Songshan Lake Materials Laboratory, Dongguan 523808, China

⁵Condensed Matter Physics and Materials Science Department,

Brookhaven National Laboratory, Upton, New York 11973, USA

(Received 19 May 2022; accepted manuscript online 16 June 2022)

As a typical hole-doped cuprate superconductor, $Bi_2Sr_2CaCu_2O_{8+\delta}(Bi2212)$ carrier doping is mostly determined by its oxygen content. Traditional doping methods can regulate its doping level within the range of hole doping. Here we report the first application of CaH₂ annealing method in regulating the doping level of Bi2212. By continuously controlling the anneal time, a series of differently doped samples can be obtained. The combined experimental results of x-ray diffraction, scanning transmission electron microscopy, resistance and Hall measurements demonstrate that the CaH₂ induced topochemical reaction can effectively change the oxygen content of Bi2212 within a very wide range, even switching from hole doping to electron doping. We also found evidence of a low-T_c superconducting phase in the electron doping side.

DOI: 10.1088/0256-307X/39/7/077403

The parent compounds of high- $T_{\rm c}$ cuprate superconductors are usually antiferromagnetic Mott insulators with the half-filling of electron occupancy. With extra electron or hole carriers, these cuprates can become superconductors with the characteristic twodome phase diagram. However, it has been noticed from the early days of the high- $T_{\rm c}$ era that the twodome phase diagram has a strong asymmetry between electron and hole doping.^[1-6] The mystery for this asymmetry, which remains unsolved to this day, is whether it is intrinsically related to fundamental differences between electron and hole doping (i.e., due to different residing orbitals^[7] or different correlation strengths caused by electron or hole charges^[8]) or the structural difference between electron-doped cuprates (T' structure) and hole-doped cuprates (Tstructure).^[9,10] To solve this mystery, it is desirable to have a cuprate material which can be doped by both electrons and holes, thus removing the influence of structure difference. Unfortunately, only very few systems can be doped by either holes or electrons.^[11-16] Previously we made efforts along this direction by developing the ozone-vacuum annealing method.^[17]

While this method can in situ dope the surface layers of hole-doped $Bi_2Sr_2CaCu_2O_{8+\delta}$ (Bi2212) and electron-doped $La_{2-x}Ce_xCuO_4$ over a wide doping range, it fails to reach the opposite doping side.^[18,19] Here we focus on the Bi2212 system again, but by using a new CaH₂ annealing method to further reduce the valence of Cu, which is inspired by the discovery of nickelate superconductors.^[20]

In this Letter, we report a systematic study of regulating the doping level of Bi2212 by the CaH₂ annealing method [Fig. 1(a)]. The original samples are optimally hole-doped Bi2212 single crystals grown by the floating-zone technique.^[21] Firstly, we select highquality optimally doped Bi2212 samples and wrap them loosely by a small piece of aluminum foil. Secondly, we take an appropriate amount of CaH₂ powder (0.18 gram are used in all treated samples) and put it into a quartz tube. Thirdly, we put the selected samples into the quartz tube without contacting CaH₂ powder and pump the tube to vacuum. Fourthly, we seal the tube after its vacuum is better than 1×10^{-4} Torr, and heat it with a tube furnace. For all the samples shown here unless mentioned oth-

[†]These authors contributed equally to this work.

^{*}Corresponding author. Email: dingh@iphy.ac.cn

 $[\]textcircled{O}$ 2022 Chinese Physical Society and IOP Publishing Ltd

erwise, we adapt the heating rate of $4.5 \,^{\circ}$ C/min until it reaches 280 $^{\circ}$ C and keep this temperature for a variable period of time (anneal time) to control the annealing sequence. Finally, we reduce the temperature to the room temperature at the same $4.5 \,^{\circ}$ C/min rate and then move the treated samples out of the tube. We select and cleave the samples with clean and flat surface for further experiments.



Fig. 1. Schematic diagram of CaH₂ annealing method and structural characterization of Bi2212 single crystals with different annealing time. (a) Schematic diagram of regulating the doping level of Bi2212 by the CaH₂ annealing method and the optical pictures of optimally doped, 4-hour annealed, and 72-hour annealed Bi2212, respectively. (b) Room-temperature XRD spectra of Bi2212 with different annealing time. All samples are annealed at 280 °C, but with different annealing time. (c) Atomically resolved iDPC-STEM image of an optimally doped Bi2212 sample. (d) Same as (c) but for a Bi2212 sample annealed for 4 h at 280 °C.

With the increase of annealing time, the luster of the sample surface disappears gradually, as shown in Fig. 1(a). In order to verify that this CaH₂ annealing method can indeed regulate the oxygen content of Bi2212 sample, we characterized the samples by x-ray diffraction (XRD) and scanning transmission electron microscopy (STEM). Figure 1(b) shows the XRD spectra of Bi2212 with different annealing time. Within a large range of annealing time (4- $36 \,\mathrm{h}$), five main diffraction peaks^[22] of (008), (0010), (0012), (0016), (0020) persist, indicating that the crystal structure is largely retained. With the increase of annealing time, the (0010) diffraction peak and the (0020) peak tend to shift to higher diffraction angles, indicating that the annealing effect reduces the c-axis length. The broadening of these two peaks may indicate the non-uniform reduction of the *c*-axis. Figure 1(c) shows the atomically resolved integrated differential phase contrast STEM (iDPC-STEM) image of an optimally doped Bi2212 sample. The interstitial oxygen atoms marked with red arrows are located between the BiO layers and SrO layers, similar to previous reports.^[23] However, after annealed for 4 h at 280 °C, the interstitial oxygen atoms disappear, as shown in Fig. 1(d). It is widely believed that the interstitial oxygen corresponds to the value of δ of the oxygen doping level,^[24,25] so this result confirms that the CaH₂ annealing method can indeed regulate the oxygen doping level of Bi2212. In addition, the atomic arrangement of the samples after CaH₂ annealing is flatter than that of the optimally doped samples, as clearly seen in Figs. 1(c) and 1(d). This means that the incommensurate modulation in Bi2212 gets weaker after CaH₂ annealing, which would lead to a decrease of the superlattice intensity.^[26]

To determine to what extent the doping level of Bi2212 system can be regulated by this CaH₂ annealing method, we perform transport measurements, as shown in Fig. 2. We measure the resistances of samples by the four-wire method^[27] and normalize them with their corresponding resistances at 300 K [Fig. 2(a)]. The resistance of the samples starts to exhibit the insulator-like behavior when annealed for 2 h, and reaches the maximum around an 8-hour annealing time, then decreases with further increase of annealing time. This behavior suggests that the samples change from the metallic/superconducting phase to the insulating phase and revert to another metallic phase. Furthermore, we measure the Hall resistance of the samples with different annealing time by the

Express Letter

Van der Pauw method.^[28] Figures 2(b) and 2(c) show the Hall resistance measured at the room temperature (300 K) and at the low temperatures (\leq 30 K), respectively. With the increase of annealing time, the sign of Hall resistance gradually changes from positive to negative, indicating that the main carrier type changes from hole to electron. To better quantify the change of resistance, we show the resistance ratios of 5 K and 300 K of the samples with different annealing time in Fig. 2(d). Apparently, the increase in resistance is because the samples approach an insulating parent phase as the hole doping level decreases. With further reduction of oxygen content, the resistance of the samples gradually decreases, which indicates that the Bi2212 system may enter the electron-doped region. Considering that the thicknesses (d) of all the samples are approximately the same, Fig. 2(e) actually shows the variation trend of Hall carrier concentration (n) with annealing time, which is calculated from Figs. 2(b) and 2(c). These transport results indicate that it is possible to change the doping type from hole to electron in the Bi2212 system by the CaH₂ annealing method.



Fig. 2. Transport results of Bi2212 samples with different annealing time. (a) Resistance curves normalized at 300 K for Bi2212 samples with different annealing time. (b) Hall resistance of the samples annealed for different time, which is measured by Van der Pauw method at the room temperature (300 K). (c) Hall resistance measured at low temperatures ($\leq 30 \text{ K}$). (d) Resistance ratio of 5 K and 300 K as a function of the annealing time. (e) Product of Hall carrier concentration (n) and sample thickness (d) versus the annealing time. To better visualization, the annealing time (2-72 h) is commuted to the logarithm form.

Next, it is desirable to find out whether the metallic state in the electron-doped side can become superconducting at a lower temperature. Therefore, we preform the lower temperature transport measurements by using a dilution refrigerator which can reach a lowest temperature of 10 mK. Remarkably, we find that a new superconducting phase appeared in some samples with a long annealing time, as shown in Fig. 3(a), which displays three samples annealed for more than 77 h at 290 °C with a much slower heating and cooling rate of $0.45 \,^{\circ}\text{C/min}$. The onset temperatures of the superconductivity $(T_{\rm c,onset})$ are ~ 1.23 K, 1.4 K and 1.44 K, respectively, and the temperatures with zero resistance $(T_{c,0})$ are ~0.86 K, 0.94 K and 1.1 K, respectively. The small deviation between $T_{c,onset}$ and $T_{\rm c,0}$ indicates the sharp superconducting transitions. The detailed superconducting properties of this new superconductor will be reported elsewhere. We also

measure the Hall resistance of the sample annealed for 84 h at 290 °C [Fig. 3(b)], and find that its carriers are indeed mainly electrons. Further transport experiments shown in the Supplementary Material (Fig. S1) strongly suggest that the superconductivity is not derived from residual hole-doped Bi2212. Furthermore, we examine the values of $T_{\rm c}$ of the elements (Bi, Sr, Ca, Cu, O) contained in the Bi2212 sample and some multi-component compounds of these elements.^[29-34] $T_{\rm c}$ of Bi bulk under the ambient pressure is below $0.53 \,\mathrm{mK}$ while $T_{\rm c}$ of cylindrical Bi nanowire is below $1.3 \,\mathrm{K}^{[35]}$ T_c of Bi element can reach 8.2 K under pressure up to 8.1 GPa while T_c of its single oxide δ -Bi₂O₃ is 5.8 K.^[36,37] As for the multi-component compounds of these elements, there are also some combinations can become high- $T_{\rm c}$ superconducting materials, such as $Bi_2Sr_2CuO_{6+\delta}$ (Bi2201) and $Bi_2Sr_2Ca_2Cu_3O_{10+\delta}$ (Bi2223). We also found trace of a low- T_c superconducting phase when the Bi2212 system is doped into the electron doping side. Considering the growth technology, pressure environment, transport properties, XRD patterns, etc., the superconductors listed above are not east to form. Nevertheless, we cannot completely exclude the possibility of trace amount. Therefore, this electron-doped superconductivity may arise from the electron-doped Bi2212 phase, or from the secondary impurity phase.



Fig. 3. Ultra-low-temperature transport results of the heavily annealed superconductors. (a) Resistance curves normalized at 2.5 K with three different annealing times of 77, 81, and 84 h, respectively. (b) Hall resistance of the sample annealed for 84 h at 290 °C. (c) A phase diagram of the Bi2212 system reached by the CaH₂ annealing method. The triangle symbols refer the $T_{c,0}$ and the green circles are the corresponding resistance values measured at 5 K in Fig. 2(a).

We summarize our results in Fig. 3(c), which roughly sketches a phase diagram that the Bi2212 samples reached by this new CaH₂ annealing process. The ability of reaching the insulting phase or even the electron side in the Bi2212 system may help us to solve the mystery of electron-hole asymmetry in the cuprates. In fact, we have conducted some preliminary ARPES results over this new phase diagram, which will be presented in a separate paper. However, one caveat for this method needs to be mentioned here: during the anneal process, hydrogen coming out of CaH₂ would combine with oxygen extracted out of Bi2212 and forms water, and the water vapor will likely degrade the Bi2212 single crystal as it comes out of the sample. To improve the single crystal quality, one can try to optimize the hydrogenation conditions or use thin film samples. The reason why we could not get higher $T_{\rm c}$ in the electron-doped samples may be that the heavily annealed samples tend to have poorer quality. In fact, the reason why electron doping is so difficult to achieve in Bi2212 and most hole-doped systems is likely due to the fact that their charge transfer gaps are larger than the electron-doped systems such as $Nd_{2-x}Ce_xCuO_4$, so a brute force method such as the CaH_2 annealing process is needed to reach the electron-doping side.

In conclusion, we develop the CaH_2 annealing method to regulate the doped carrier of Bi2212 single crystal samples over an unprecedently wide doping range, and may even reach the superconducting electron-doped phase. This will likely open a new venue for systematically studying the cuprate systems from both electron and hole sides, from which new results would provide new insights on some of the most fundamental issues the high- $T_{\rm c}$ cuprate superconductors are facing, including the electron-hole asymmetry, the pseudogap, and the high- $T_{\rm c}$ superconducting mechanism.

Acknowledgments. We thank Xing-Chen Guo, Zhe Zheng, Feng Ran and Ren-Jie Zhang for technical support, and Zheng-Yu Weng, Fu-Chun Zhang and Tao Xiang for helpful discussions. This work was supported by the National Natural Science Foundation of China (Grant Nos. 11888101 and U1832202), the Chinese Academy of Sciences (Grant Nos. QYZDB-SSW-SLH043 and XDB33000000), the K. C. Wong Education Foundation (Grant No. GJTD-2018-01), and the Informatization Plan of Chinese Academy of Sciences (Grant No. CAS-WX2021SF-0102). This work was also supported by the Synergetic Extreme Condition User Facility (SECUF). Y.L.G. was supported by China Postdoctoral Science Foundation (Grant Nos. 2020M680726 and YJ20200325). G.D.G. was supported by US DOE (Grant Nos. DE-SC0010526 and DE-SC0012704).

References

- [1] Bednorz J G and Müller K A 1986 Z. Phys. B 64 189
- [2] Tokura Y, Takagi H and Uchida S 1989 Nature 337 345
- [3] Jang S W, Sakakibara H, Kino H, Kotani T, Kuroki K and Han M J 2016 Sci. Rep. 6 33397
- [4] Rybicki D, Jurkutat M, Reichardt S, Kapusta C and Haase J 2016 Nat. Commun. 7 11413
- [5] Jin C Q, Laffez P, Tatsuki T, Tamura T, Adachi S, Koshizuka N Y H, Tanaka S and Wu X J 1995 Nature 375 301

- [6] Mazumdar S 2018 Phys. Rev. B 98 205153
- [7] Naito M, Krockenberger Y, Ikeda A and Yamamoto H 2016 Physica C 523 28
- [8] Weber C, Haule K and Kotliar G 2010 Nat. Phys. 6 574
- [9] Li Y M, Tabis W, Tang Y, Yu G, Jaroszynski J, Barišić N and Greven M 2019 Sci. Adv. 5 eaap7349
- [10] Hirsch J E and Marsiglio F 2019 Physica C 564 29
- [11] Segawa K, Kofu M, Lee S H, Tsukada I, Hiraka H, Fujita M, Chang S, Yamada K and Ando Y 2010 Nat. Phys. 6 579
 [12] Adachi T, Kawamata T and Koike Y 2017 Condens. Matter
- [13] Hu C, Zhao J F, Gao Q, Yan H T, Rong H T, Huang J W, Liu J, C, Y Q, Li C, Chen H, Zhao L, Liu G D, Jin C Q, Xu Z Y, Xiang T and Zhou X J 2021 Nat. Commun. 12 1356
- [14] Segawa K and Ando Y 2006 Phys. Rev. B 74 100508
- [15] Zhong Y, Fan J Q, Wang R F, Wang S Z, Zhang X F, Zhu Y Y, Dou Z Y, Yu X Q, Wang Y, Zhang D, Zhu J, Song C L, Ma X C and Xue Q K 2020 *Phys. Rev. Lett.* **125** 077002
- [16] Zeng S W, Wang X, Lü W M, Huang Z, Motapothula M, Liu Z Q, Zhao Y L, Annadi A, Dhar S, Mao H, Chen W, Venkatesan T and A 2012 *Phys. Rev. B* 86 045124
- [17] Zhong Y G, Guan J Y, Shi X, Zhao J, Rao Z C, Tang C Y, Liu H J, Weng Z Y, Wang Z Q, Gu G D, Qian T, Sun Y J and Ding H 2018 *Phys. Rev. B* 98 140507
- [18] Zhong Y G, Guan J Y, Zhao J, Tang C Y, Rao Z C, Liu H J, Zhang J H, Li S, Weng Z Y, Gu G D, Sun Y J and Ding H 2019 Phys. Rev. B 100 184504
- [19] Tang C Y, Lin Z F, Zhang J X, Guo X C, Guan J Y, Gao S Y, Rao Z C, Zhao J, Huang Y B, Qian T, Weng Z Y, Jin K, Sun Y J and Ding H 2021 *Phys. Rev. B* 104 155125
- [20] Li D F, Lee K, Wang B Y, Osada M, Crossley S, Lee H R, Cui Y, Hikita Y and Hwang H Y 2019 Nature 572 624
- [21] Wen J S, Xu Z J, Xu G Y, Hücker M, Tranquada J M and Gu G D 2008 J. Cryst. Growth 310 1401
- [22] Jindal A, Jangade D A, Kumar N, Vaidya J, Das I, Bapat R, Parmar J, Chalke B A, Thamizhavel A and Deshmukh M M 2017 Sci. Rep. 7 3295

- [23] Song D S, Zhang X F, Lian C S L H, Alexandrou I, Lazić I, Bosch E G T, Zhang D, Wang L L, Yu R, Cheng Z Y, Song C L, Ma X C, Duan W H, Xue Q K and Zhu J 2019 Adv. Funct. Mater. 29 1903843
- [24] Presland M R, Tallon J L, Buckley R G, Liu R S and Flower N E 1991 Physica C ${\bf 176}$ 95
- [25] Tallon J L, Bernhard C, Shaked H, Hitterman R L and Jorgensen J D 1995 Phys. Rev. B 51 12911
- [26] Liu J, Zhao L, Gao Q et al. 2019 Chin. Phys. B 28 077403
- [27] Quaranta Q, Gades L M, Xue C, Divan R, Patel U M, Guruswamy T and Miceli A 2021 arXiv:2111.02503 [condmat.mtrl-sci]
- [28] Thiery N, Naletov V V, Vila L, Marty A, Brenac A, Jacquot J F, Loubens G D, Viret M, Anane A, Cros V, Youssef J B, Beaulieu N, Demidov V E, Divinskiy B, Demokritov S O and Klein O 2018 Phys. Rev. B 97 064422
- [29] Prakash O, Kumar A, Thamizhavel A and Ramakrishnan S 2017 Science 355 52
- [30] Medina J C, Bizarro M, Silva-Bermudez P, Giorcelli M, Tagliaferro A and Rodil S E 2016 *Thin Solid Films* 612 72
- [31] Thaowonkaew S, Chao-moo W, Nontra-udorn R, Vora-ud A and Seetawan T 2017 Mater. Today Proc. 4 6592
- [32] Piriou A, Giannini E, Fasano Y, Senatore C and Fischer Y 2010 Phys. Rev. B 81 144517
- [33] Adachi S, Usui T, Takahashi K, Kosugi K, Watanabe T, Nishizaki T, Adachi T, Kimura S, Sato K, Suzuki K M, Fujita M, Yamada K and Fujii T 2015 *Phys. Proceedia* 65 53
- [34] Altın S, Aksan M A, YakıN M E and BalcıY 2010 J. Alloys Compd. 502 16
- [35] Tian M L, Wang J, Ning W, Mallouk T E and Chan M H W 2015 Nano Lett. 15 1487
- [36] Li Y F, Wang E Y, Zhu X Y and Wen H H 2017 *Phys. Rev.* B 95 024510
- [37] Koza J A, Bohannan E W and Switzer J A 2013 Acs Nano 7 9940