

Time-Reversal Symmetry Breaking Driven Topological Phase Transition in EuB_6

Shun-Ye Gao^{1,2,*}, Sheng Xu^{3,4,*}, Hang Li^{5,6,*}, Chang-Jiang Yi^{1,*}, Si-Min Nie⁷, Zhi-Cheng Rao^{1,2}, Huan Wang^{3,4},
 Quan-Xin Hu^{1,2}, Xue-Zhi Chen^{2,8}, Wen-Hui Fan^{1,2}, Jie-Rui Huang^{1,2}, Yao-Bo Huang⁹, Nini Pryds⁵, Ming Shi⁶,
 Zhi-Jun Wang^{1,2}, You-Guo Shi^{1,2,10,§}, Tian-Long Xia^{3,4,‡}, Tian Qian^{1,10,†} and Hong Ding^{1,2,10,11}

¹Beijing National Laboratory for Condensed Matter Physics and Institute of Physics,
 Chinese Academy of Sciences, Beijing 100190, China

²University of Chinese Academy of Sciences, Beijing 100049, China

³Department of physics, Renmin University of China, Beijing 100872, China

⁴Beijing Key Laboratory of Opto-electronic Functional Materials & Micro-nano Devices,
 Renmin University of China, Beijing 100872, China

⁵Department of Energy Conversion and Storage, Technical University of Denmark,
 Fysikvej Building 310, DK-2800 Kgs. Lyngby, Denmark

⁶Swiss Light Source, Paul Scherrer Institut, CH-5232 Villigen, Switzerland

⁷Department of Materials Science and Engineering, Stanford University, Stanford, California 94305, USA

⁸Shanghai Institute of Applied Physics, Chinese Academy of Sciences, Shanghai 201800, China

⁹Shanghai Synchrotron Radiation Facility, Shanghai Advanced Research Institute,
 Chinese Academy of Sciences, Shanghai 201204, China

¹⁰Songshan Lake Materials Laboratory, Dongguan, Guangdong 523808, China

¹¹CAS Center for Excellence in Topological Quantum Computation,
 University of Chinese Academy of Sciences, Beijing 100049, China



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The interplay between time-reversal symmetry (TRS) and band topology plays a crucial role in topological states of quantum matter. In time-reversal-invariant (TRI) systems, the inversion of spin-degenerate bands with opposite parity leads to nontrivial topological states, such as topological insulators and Dirac semimetals. When the TRS is broken, the exchange field induces spin splitting of the bands. The inversion of a pair of spin-splitting subbands can generate more exotic topological states, such as quantum anomalous Hall insulators and magnetic Weyl semimetals. So far, such topological phase transitions driven by the TRS breaking have not been visualized. In this work, using angle-resolved photoemission spectroscopy, we have demonstrated that the TRS breaking induces a band inversion of a pair of spin-splitting subbands at the TRI points of Brillouin zone in EuB_6 , when a long-range ferromagnetic order is developed. The dramatic changes in the electronic structure result in a topological phase transition from a TRI ordinary insulator state to a TRS-broken topological semimetal (TSM) state. Remarkably, the magnetic TSM state has an ideal electronic structure, in which the band crossings are located at the Fermi level without any interference from other bands. Our findings not only reveal the topological phase transition driven by the TRS breaking, but also provide an excellent platform to explore novel physical behavior in the magnetic topological states of quantum matter.

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*These authors contributed equally to this work.

†Corresponding author.

tqian@iphy.ac.cn

‡Corresponding author.

tlxia@ruc.edu.cn

§Corresponding author.

ygshe@iphy.ac.cn

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I. INTRODUCTION

Since the discovery of the integer quantum Hall effect, time-reversal symmetry (TRS) has been a key issue in the study of topological states of quantum matter [1–3]. The thinking and research on TRS directly led to the discovery of time-reversal-invariant (TRI) topological insulators with the quantum spin Hall effect [4–9]. In magnetic materials, the TRS breaking induced by magnetic order can produce more exotic topological states. The most famous examples are the realization of quantum anomalous Hall (QAH) effect [10–19] and the discovery of magnetic Weyl

semimetals [20–23]. The interplay of magnetism and band topology may lead to the emergence of novel physical behavior, which would give new access to applications of topological states of quantum matter.

The QAH insulator and magnetic Weyl semimetal states are closely related to the band inversion under the TRS breaking. As illustrated in Fig. 1(a), regardless of whether the initial TRI systems with inversion symmetry are

topological insulators with band inversion or ordinary insulators, when the TRS is broken, the exchange field leads to an inversion of one pair of spin-splitting subbands with opposite chirality at the TRI point, whereas the other pair is separated. In two-dimensional (2D) systems, the band inversion opens an energy gap characterized by a nonzero Chern number, forming the QAH insulators [10]. In three-dimensional systems, the band inversion cannot be

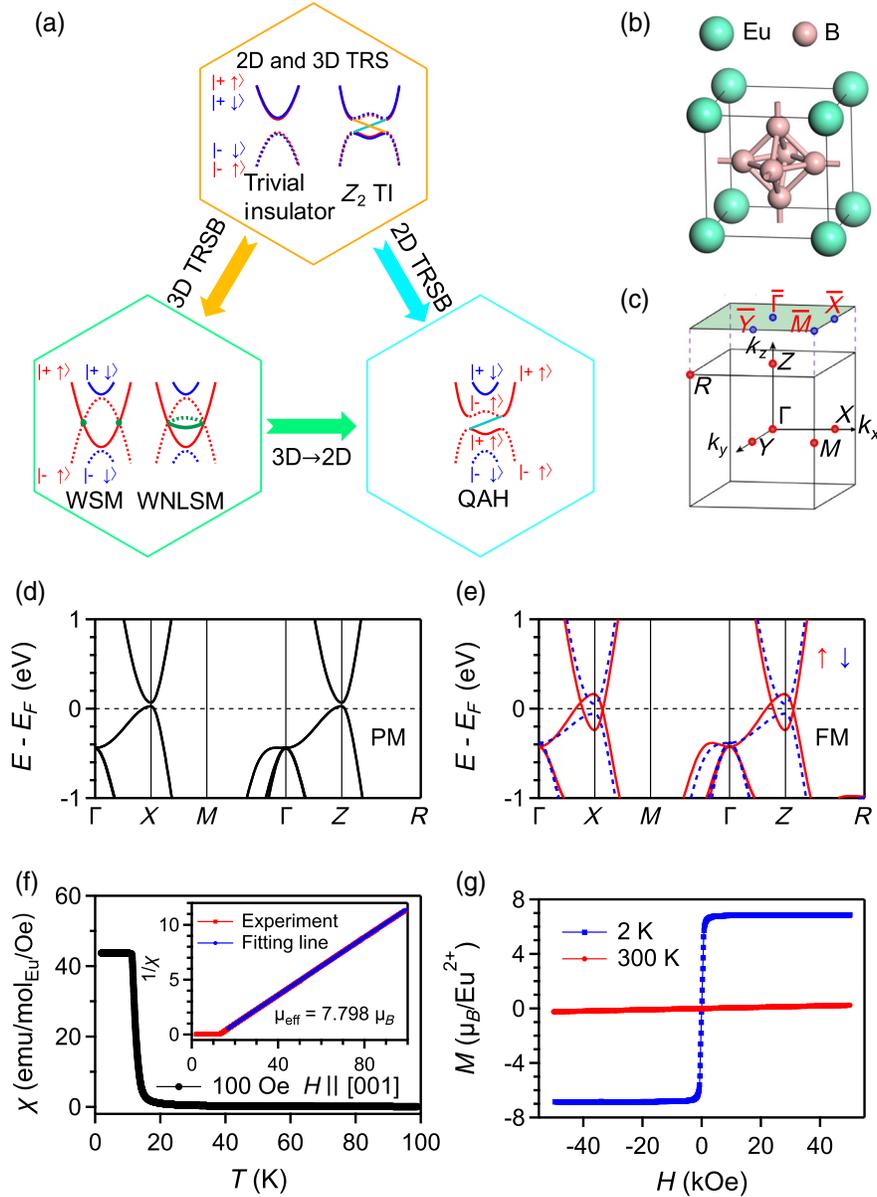


FIG. 1. Calculated electronic structures and magnetic properties of EuB_6 . (a) Schematic plots of emergent topological phases with TRS breaking (TRSB) in centrosymmetric 3D and 2D systems. Topological insulator (TI), Weyl semimetal (WSM), and Weyl nodal-line semimetal (WNLSM). “+” and “−” represent opposite parity of the subbands. “ \uparrow ” and “ \downarrow ” represent spin up and spin down. (b) Crystal structure of EuB_6 . (c) Bulk BZ and (001) surface BZ. (d), (e) Calculated electronic structures along high-symmetry lines in the PM and FM states, respectively. Magnetic moments are defined to be along the [001] direction in the FM state. The calculated results are adopted from Ref. [24]. (f) Temperature dependence of magnetic susceptibility. Inset: temperature dependence of inverse susceptibility. The blue line is a fitting to the Curie-Weiss formula. (g) Magnetic field dependence of magnetization at 2 K (blue) and 300 K (red) with magnetic fields parallel to the [001] direction.

fully gapped even with spin-orbit coupling, forming the magnetic topological semimetals (TSMs) with twofold degenerate Weyl points or nodal lines [24]. Under the 2D limit, the degeneracy is lifted by opening an energy gap and the magnetic TSMs are transformed into the QAH insulators [20,24].

While it is believed that the QAH insulator state in magnetically doped Bi_2Te_3 [10–15] and MnBi_2Te_4 [16–19] and the magnetic Weyl semimetal state in $\text{Co}_3\text{Sn}_2\text{S}_2$ [20–23] arise from the mechanism illustrated in Fig. 1(a), the topological phase transitions driven by the TRS breaking have not been observed. In this work, we use angle-resolved photoemission spectroscopy (ARPES) to visualize the topological phase transition from a TRI ordinary insulator state to a TRS-broken TSM state in EuB_6 , which is accompanied by the development of a long-range ferromagnetic (FM) order. Remarkably, the band crossings are the only feature at the Fermi level (E_F) without the existence of other Fermi surfaces (FSs), which is favorable to the emergence of exotic physics associated with the topological nodes.

II. METHODS

Single crystals of EuB_6 were grown by the flux method. The ingredients with a ratio of $\text{Eu}:\text{B}:\text{Al} = 1:6:570$ were put into the alumina crucible with a cap and sealed into a quartz ampoule. Then the ampoule was heated to 1350°C , kept for 10 h, and cooled to 700°C at a rate of 5°C per hour. EuB_6 single crystals were obtained after the Al flux was removed by dissolving in sodium hydroxide solution. All the ARPES data shown were collected at the “dreamline” beam line and the 03U beam line at the Shanghai Synchrotron Radiation Facility, the SIS-HRPES beam line at the Swiss Light Source, and the BL25SU beam line at the Spring 8. We conducted complementary experiments at the 13U beam line at the National Synchrotron Radiation Laboratory at Hefei, the “Bloch” beam line, MAXIV, Sweden, the 7.0.2.1 beam line, ALS, U.S., and the I05 beam line, Diamond, UK. The energy and angular resolutions were set to 15–70 meV and 0.1° . The samples were cleaved *in situ* and measured in ultrahigh vacuum better than 7×10^{-11} torr.

III. RESULTS

EuB_6 has the CsCl-type crystal structure with space group $Pm\bar{3}m$ (no. 221) [25]. As seen in Fig. 1(b), Eu atoms are located at the corners of cubic lattice and six B atoms form an octahedral cage inside the lattice. The magnetic susceptibility (χ) curves in Fig. 1(f) show that EuB_6 is in the paramagnetic (PM) state at high temperatures. The fitting to the Curie-Weiss formula gives an effective magnetic moment $\mu_{\text{eff}} = 7.798\mu_B/\text{Eu}^{2+}$, which is close to the theoretical value $7.94\mu_B$ of local $4f^7$ orbitals. The magnetic susceptibility starts to rise sharply at ~ 15 K and reaches a

plateau at ~ 11 K in Fig. 1(f), indicating a PM to FM transition consistent with previous reports [26–28]. Specific heat measurements revealed two consecutive phase transitions at $T_M = 15.5$ K and $T_C = 12.6$ K [29], which were subsequently confirmed and extensively studied by various experimental techniques [30–33]. These studies indicate that isolated magnetic polarons start to form at ~ 30 K, and they grow in size and/or increase in number with decreasing temperature. The magnetic polarons overlap and percolate in the bulk at T_M , forming a ferromagnetically ordered phase coexisting with the PM phase. With further decreasing temperature, the volume fraction of the FM phase increases at the expense of the PM phase, and eventually a complete FM state is achieved at T_C . The adiabatic magnetization curve at 2 K in Fig. 1(g) shows that EuB_6 is a soft ferromagnet with an extremely low coercive field ~ 35 Oe.

Theoretical calculations indicate that EuB_6 is a semiconductor with trivial topology in the PM state [24]. Both valence band top and conduction band bottom are located at three TRI points, X , Y , and Z , which are equivalent in the PM state, forming a direct band gap of ~ 20 meV at E_F in Fig. 1(d). The valence and conduction bands near E_F are mainly from the B $2p$ and Eu $5d$ states, respectively. When the Eu^{2+} $4f$ magnetic moments form a FM order, the interatomic antiferromagnetic exchange coupling between the B $2p$ and occupied (unoccupied) Eu $4f$ states pushes the spin-up (spin-down) valence subband upward (downward), while the intra-atomic FM exchange coupling between the Eu $5d$ and occupied (unoccupied) $4f$ states pushes the spin-up (spin-down) conduction subbands downward (upward) [24,34,35]. Because of the special exchange coupling, the band inversion occurs at three TRI points in the spin-up channel, while the band gap increases in the spin-down channel, as seen in Fig. 1(e). According to theoretical analysis [24], the band inversion at these TRI points cannot be fully gapped, forming Weyl points or nodal lines with twofold degeneracy, depending on the magnetic moment direction. When the magnetic moments are aligned with the $[001]$ direction, the mirror symmetry \hat{M}_z protects three nodal lines encircling three TRI points, respectively. When the magnetic moments are aligned with the $[110]$ direction, the mirror symmetry \hat{M}_{110} protects one nodal line encircling the Z point while Weyl points can be found near the X and Y points. When the magnetic moments deviate from the $[001]$ and $[110]$ directions, no nodal lines survive because the mirror symmetries are broken, but Weyl points always exist near three TRI points and their positions move with varying the magnetic moment direction.

We first measure electronic structures in the PM state ($T = 20$ K) on (001) cleavage surfaces of EuB_6 with ARPES. We observe two sets of distinct electronic structures near E_F , and the corresponding core level spectra have different peak structures. As seen in Fig. 2(a), the

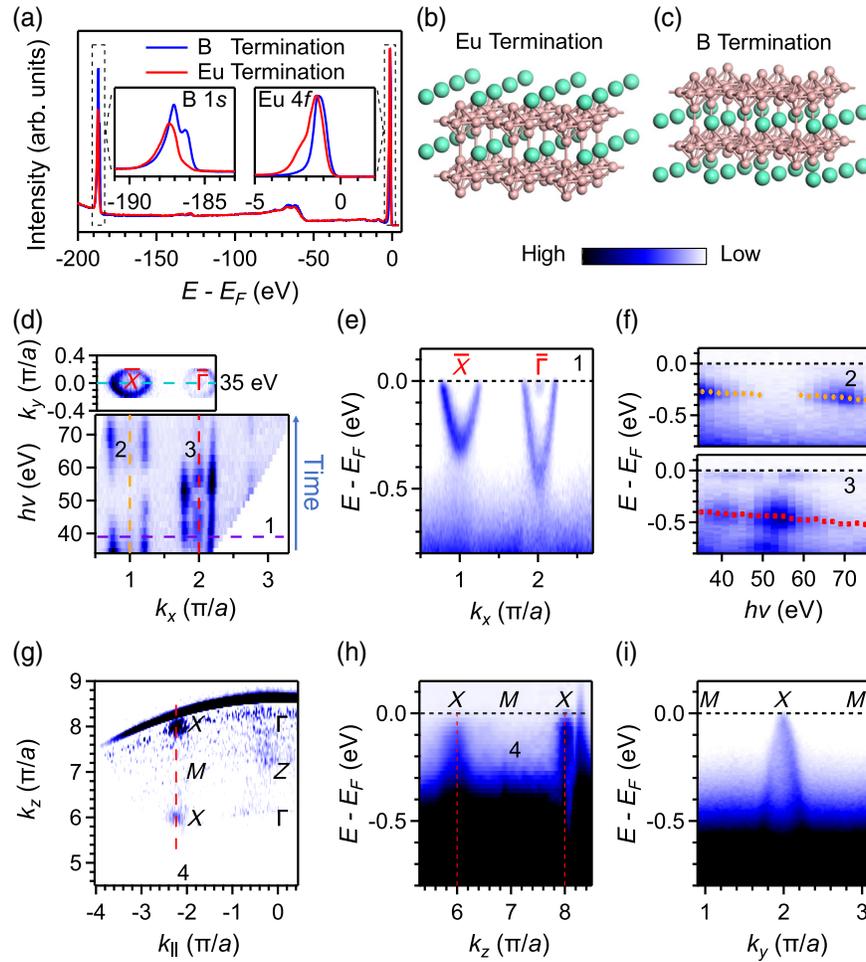


FIG. 2. Electronic structures measured on Eu- and B-terminated surfaces in the PM state. (a) Core-level photoemission spectra showing characteristic peaks of the B $1s$ and Eu $4f$ orbitals in the insets. (b),(c) Schematic plots of (001) cleavage surfaces with Eu and B terminations, respectively. (d) Top: intensity plot of the ARPES data at E_F with $h\nu = 35$ eV. Bottom: intensity plot of the ARPES data at E_F collected in a range of photon energies from 34 to 76 eV. (e) Intensity plot of the ARPES data along cut 1. (f) Intensity plots of the ARPES data along cuts 2 and 3, showing photon energy and time dependence of the band bottoms. (g) Intensity plot of the ARPES data at E_F collected along the direction from $(0,0)$ to $(-\pi/a, 2\pi/a)$ in a range of photon energies from 59 to 151 eV. The inner potential $V_0 = 17.73$ eV and the lattice constant $a = 4.187$ Å. (h) Intensity plot of the ARPES data along cut 4. (i) Intensity plot of the ARPES data along M - X in the k_x - k_y plane with $k_z = 8\pi/a$. The data in (d)–(f) and (g)–(i) were collected on Eu- and B-terminated surfaces, respectively. All the data were collected at 20 K.

B $1s$ (Eu $4f$) states have single-peak (double-peak) and double-peak (single-peak) structures, respectively. As chemical environments of the atoms in the outmost layer are different from those of the atoms underneath, the atoms in the outmost layer can produce additional peaks in the core level states. We attribute the double-peak structures of the B $1s$ and Eu $4f$ states to two kinds of terminated surfaces with either B or Eu atoms in the outmost layer, as illustrated in Figs. 2(b) and 2(c). Recent scanning tunnelling microscopy measurements on (001) cleavage surfaces of EuB_6 found two kinds of distinctive topographies, which were assigned to Eu- and B-terminated surfaces, respectively [36,37]. By scanning the Eu $4f$ and B $1s$ core levels at different positions on the cleavage surfaces with incident light, whose spot size varies between 30×50 and $50 \times 150 \mu\text{m}^2$

at various synchrotron beam lines, we can find single termination areas with only one set of electronic structures. The results were repeated in independent measurements on eight samples.

The electronic structures measured on Eu-terminated surfaces are plotted in Figs. 2(d)–2(f). We observe two electron bands near E_F at $\bar{\Gamma}$ and \bar{X} , respectively, where a is the lattice constant. Figure 2(f) shows that the bottoms of both bands have no obvious dispersion with varying photon energy ($h\nu$) except for the downward shift during measurements (see more details on time-dependent changes of the bands in Fig. 5 and 6 in the Appendix A). The nondispersive characteristic along the surface-perpendicular momentum direction k_z indicates that the bands observed on Eu-terminated surfaces are surface

states. The electron band at the \bar{X} point as well as its energy shift with time was also found in previous ARPES experiments [38,39], where it was speculated that the freshly cleaved surfaces could be metastable and relax with time, resulting in time-dependent band bending.

On the contrary, the $h\nu$ -dependent data collected on B-terminated surfaces in Figs. 2(g) and 2(h) exhibit an obvious k_z dispersion with a period of $2\pi/a$, which is an indication of bulk states. We observe near- E_F electronic states around the X point, which show a hole band with top very close to E_F in Fig. 2(i). Moreover, a shallow electron band just above the hole band is observed in some samples (see Fig. 7 in Appendix B). The observation is well consistent with the calculated bulk electronic structure of the PM state in Fig. 1(d), and it was not captured in previous ARPES experiments [38,39]. It is expected that the near- E_F electronic states can be observed at $\bar{\Gamma}$ and \bar{X} (or \bar{Y}), which are the (001)-surface projections of three TRI points. However, we find that the hole band is only clear at $(\pi/a, 2\pi/a)$ on the second BZ boundary, while it is unclear or almost invisible at other locations, such as $(0, 0)$, $(\pi/a, 0)$, and $(2\pi/a, 0)$.

When the samples are cooled down from the PM to FM state, the surface states on Eu-terminated surfaces do not change except for the downward shift with time (Fig. 6

in Appendix A), whereas the near- E_F electronic states on B-terminated surfaces exhibit dramatic changes. In Fig. 3(a) one can see clear band splitting with decreasing temperature through T_C on B-terminated surfaces. After the temperature-dependent measurements, we heat up the sample to 17 K and confirm that the band dispersions can be restored to the initial state in the PM state, thus excluding the possibility of time-dependent changes. The band splitting is an indication of symmetry breaking, such as the breaking of TRS or inversion symmetry. The former induces the Zeeman splitting with energy shifts while the latter induces the Rashba splitting with momentum shifts. In EuB_6 , it is clear that the TRS is broken at low temperatures due to the long-range FM order, whereas x-ray and neutron diffraction experiments did not find any signature of a structural phase transition around T_C [30,40], indicating that the inversion symmetry is preserved in the FM state. Therefore, we attribute the dramatic changes to the Zeeman splitting, which leads to energy shifts of the bands.

As seen in Fig. 3(b), there is a hole band (α) below E_F at the X point at 20 K, and the band splits into two subbands (α_1 and α_2) at 5 K. The two subbands move up and down by ~ 200 meV, respectively, from 20 to 5 K. On the other hand, we observe that an electron band (β_1) moves down with decreasing temperature. Because of the TRS breaking,

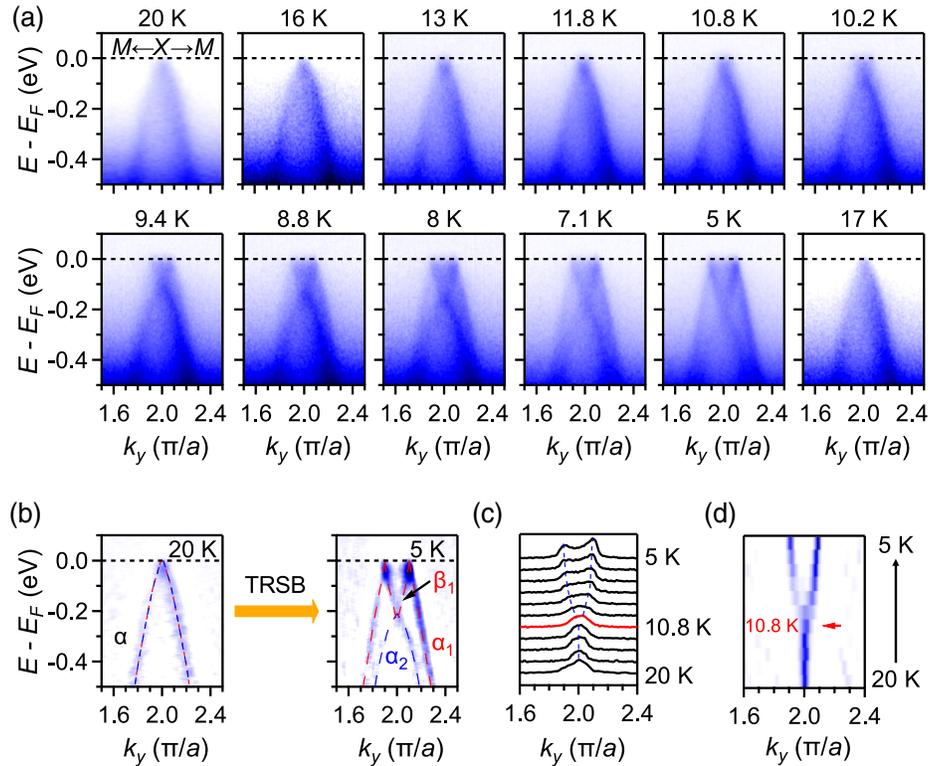


FIG. 3. Temperature dependence of band dispersions on B-terminated surfaces. (a) Intensity plots of the ARPES data along M - X at different temperatures between 20 and 5 K. (b) 2D curvature intensity plots of the data in (a) at 20 and 5 K. Dashed lines are guide to eyes of the band dispersions. (c) MDCs at E_F of the ARPES data in (a). (d) Curvature intensity plot of the data in (c). All the data were collected with $h\nu = 135$ eV.

the conduction band just above E_F should also split into two subbands. One of them moves down across E_F , whereas the other should move up and therefore cannot be detected in the ARPES experiments. The dramatic changes in the electronic structure through T_C are consistent with optical conductivity experiments, which

revealed a substantial shift toward higher frequencies of the plasma edge below T_C , suggesting an increase of itinerant charge carriers and/or a reduction of their effective mass [29].

Notably, the band splitting results in the band inversion of two subbands α_1 and β_1 at the X point. The band

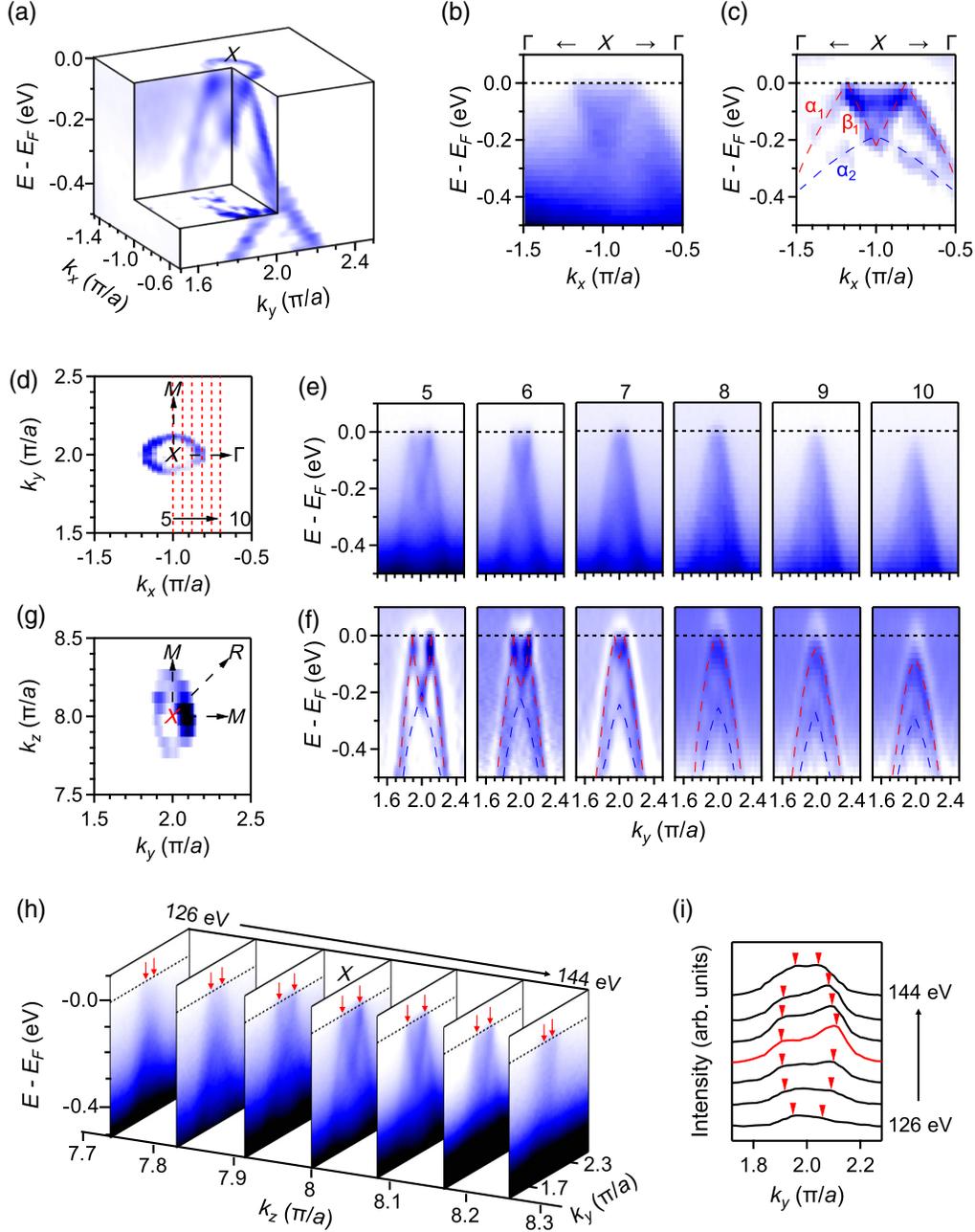


FIG. 4. In-plane and out-of-plane band dispersions on B-terminated surfaces in the FM state. (a) 3D curvature intensity plot of the ARPES data around the X point. (b) Intensity plot of the ARPES data along Γ - X . (c) 2D curvature intensity plot of the data in (b). Dashed lines are guide to eyes of the band dispersions. (d) Curvature intensity plot of the ARPES data at E_F around the X point in the k_x - k_y plane. (e) Intensity plots of the ARPES data along cuts 5–10. (f) Curvature intensity plots of the ARPES data in (e). (g) Curvature intensity plot of the ARPES data at E_F around the X point in the k_y - k_z plane. (h) Intensity plots of the ARPES data along M - X collected in a range of photon energies from 126 to 144 eV. (i) MDCs at E_F of the ARPES data in (h). The data in (a)–(f) were collected with $h\nu = 135$ eV. All the data were collected at 4.5 K.

inversion agrees well with the calculations in the FM state in Fig. 1(e) [24]. Moreover, the calculations indicate that two subbands α_2 and β_1 are also inverted in Fig. 1(e), while the experimental data show that they are almost degenerate at ~ 0.2 eV below E_F in Fig. 3(b). In the experimental data both α_2 and β_1 move down significantly, whereas in the calculations only β_1 moves down significantly while α_2 moves down slightly, suggesting that the exchange coupling between the B $2p$ and unoccupied Eu $4f$ bands is underestimated in the calculations [24]. Despite the discrepancy, both experiment and calculation consistently confirm the band inversion of α_1 and β_1 , which is critical for the magnetic TSM state in EuB_6 .

Figures 3(c) and 3(d) show temperature-dependent momentum distribution curves (MDCs) at E_F from the PM to FM state. The MDCs exhibit a single peak at 20 K, corresponding to the hole band top in the PM state. The peak starts to split at ~ 11 K, which can be regarded as an indication of the band inversion. We notice that the band dispersions are not well defined between 11 and 8 K in Fig. 3(a). The phenomenon is probably related to the phase separation between T_M and T_C , where the FM phase formed by magnetic polarons coexists with the PM phase [29–33]. The band dispersions become clear below 8 K, probably because the complete FM state is achieved. The temperature range between 11 and 8 K is lower than the two transition temperatures $T_M = 15.5$ K and $T_C = 12.6$ K determined by specific heat and magnetization measurements [29,30]. In our ARPES experiments, the actual sample temperature may be higher than that detected by the sensor because the samples are exposed to thermal radiation from the room-temperature chamber and analyzer.

We systematically investigate in-plane and out-of-plane band dispersions near the X point in the FM state. The results in Figs. 4(a)–4(f) exhibit that the band crossings at E_F form a nodal-ring FS encircling the X point in the k_x - k_y plane. As seen in Fig. 4(d), the FS is elliptical with its long axis along Γ - X . Similarly, the $h\nu$ -dependent data in Figs. 4(g)–4(i) show a nodal-ring FS formed by the band crossings in the k_y - k_z plane. The FS should be nearly isotropic with fourfold symmetry in the M - X - R plane, while the experimental FS in Fig. 4(g) appears to be stretched along k_z , which is caused by momentum broadening in the direction perpendicular to the sample surface. The FS formed by the band crossings can be regarded as the degeneracy of hole and electron FSs. de Haas–van Alphen and Shubnikov–de Haas (SdH) experiments found both hole and electron FSs with different sizes in the FM state [41,42], which could arise from the FS splitting when the chemical potential slightly deviates from the band crossings due to small nonstoichiometry. The extremal areas of the ellipsoidal FS in the ARPES data are estimated to be 278 and 410 T in the X - M - R and Γ - X - M planes, respectively, which are $\sim 15\%$ smaller than those determined by the SdH experiments at 0.4 K [42]. As seen in Fig. 3(d), the Fermi

wave vector continues to increase below 11 K and does not reach saturation at 5 K. The discrepancy in FS size can be attributed to the temperature difference between ARPES and SdH experiments.

Finally, we discuss the effects of magnetic domains. As magnetic fields are not allowed in the ARPES experiments, magnetic domains must exist in the soft ferromagnet EuB_6 . Because of the small magnetic anisotropy in EuB_6 , the typical domain size at zero field was estimated to be on the order of $1 \mu\text{m}$ [30], which is much smaller than the spot size of incident light, suggesting that the ARPES results should be a mixture of electronic structures of multiple domains with various magnetic moment directions. As the magnetic moment direction determines the band crossings at E_F , specific topological states cannot be distinguished in the mixed electronic structure of multiple domains. On the other hand, the effects of magnetic moment direction on the overall band dispersions are negligible [24]. Therefore, despite the existence of magnetic domains, clear band dispersions can be observed when the samples enter the complete FM state at low temperatures.

IV. OUTLOOK

Our ARPES data have demonstrated that the TRS breaking results in a topological phase transition with the band inversion at the X point in EuB_6 , which is well consistent with the band calculations [24]. According to theoretical analysis, the band inversion at the TRI point cannot be fully gapped, forming Weyl points or nodal lines depending on the magnetic moment direction. As a soft ferromagnet, it is easy to align the magnetic moments with external magnetic fields, which would be favorable to tune the topological states and relevant physical properties. Compared with the other experimentally discovered magnetic TSMs [21,23,43], the distinguishing characteristic in EuB_6 is that the band crossings are the only feature at E_F without the existence of other FSs, and therefore their energy positions are restricted at E_F by electron-hole compensation in stoichiometric samples.

Band calculations indicate that the band inversion occurs in the spin-up channel whereas the band gap remains in the spin-down channel, suggesting that EuB_6 is also a half-metal with a full spin polarization at E_F in the FM state [24,33,44]. Our experimental data are consistent with the predicted half-metallic electronic structure. However, the expected full spin polarization appears to be in conflict with Andreev reflection spectroscopy results showing a spin polarization of $\sim 50\%$, which is attributed to the scenario that one band is fully spin polarized while the other is unpolarized [45]. It has been argued that the half-metallic state can be tuned into an incomplete spin-polarized state by slight doping [34]. Further studies are required to resolve the disagreement.

Since Weyl points can only be defined in 3D momentum space, they will be removed by opening an energy gap

under the 2D limit. If the band order is maintained during the crossover from 3D to 2D, the Weyl semimetal state can be transformed into the QAH insulator state. Band calculations have predicted that the QAH insulator state can be achieved under the 2D limit in EuB_6 , where the band crossings open an energy gap [24]. As there is no interference from other bands, the QAH insulator state has a global gap at E_F in the density of states. These outstanding properties make EuB_6 a promising system to study the physical behavior associated with a range of magnetic topological states in both 3D crystals and 2D thin films.

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APPENDIX A: TIME DEPENDENCE OF BAND DISPERSIONS MEASURED ON EU-TERMINATED SURFACES

We study the changes of the electronic structures with time after cleavage on Eu-terminated surfaces, where two electron bands are observed near E_F at $\bar{\Gamma}$ and \bar{X} , respectively. The bands shift monotonically downwards with time at a constant temperature $T = 60$ K (Fig. 5) and under temperature cycling (Fig. 6). We note that the results in Fig. 6 show that the FM transition has negligible effects on the band dispersions on Eu-terminated surfaces, which is significantly different from dramatic changes through the FM transition on B-terminated surfaces in Fig. 3(a). This behavior implies that the $4f$ magnetic moments in the outmost layer on Eu-terminated surfaces may be disordered at low temperatures, which needs further study.

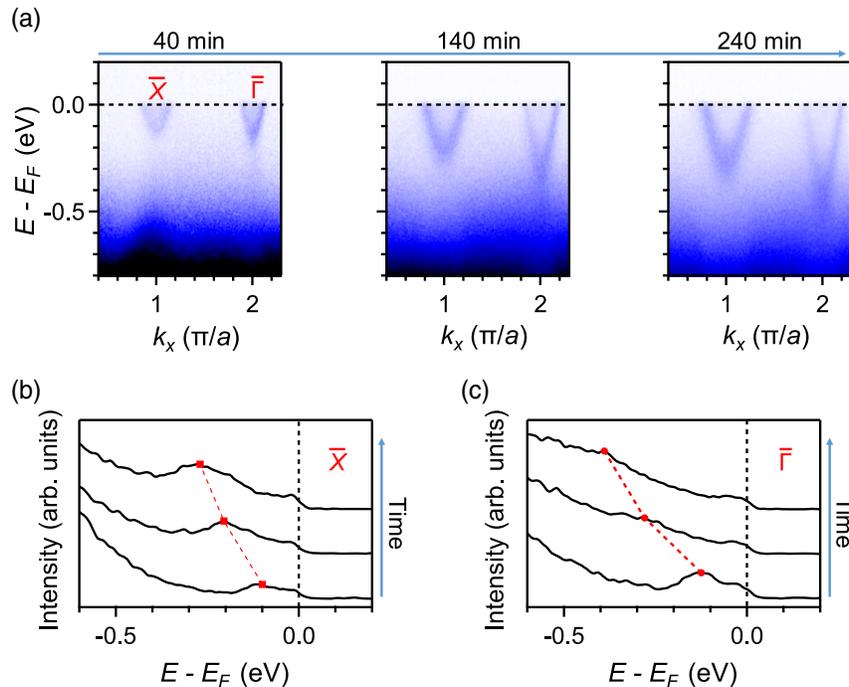


FIG. 5. Time dependence of band dispersions measured on Eu-terminated surfaces at 60 K. (a) Intensity plots of the ARPES data collected at 40, 140, and 240 min after the cleavage with $h\nu = 38$ eV. (b),(c) Energy distribution curves at the \bar{X} and $\bar{\Gamma}$ points, respectively, showing energy shifts with time of the band bottoms at 60 K.

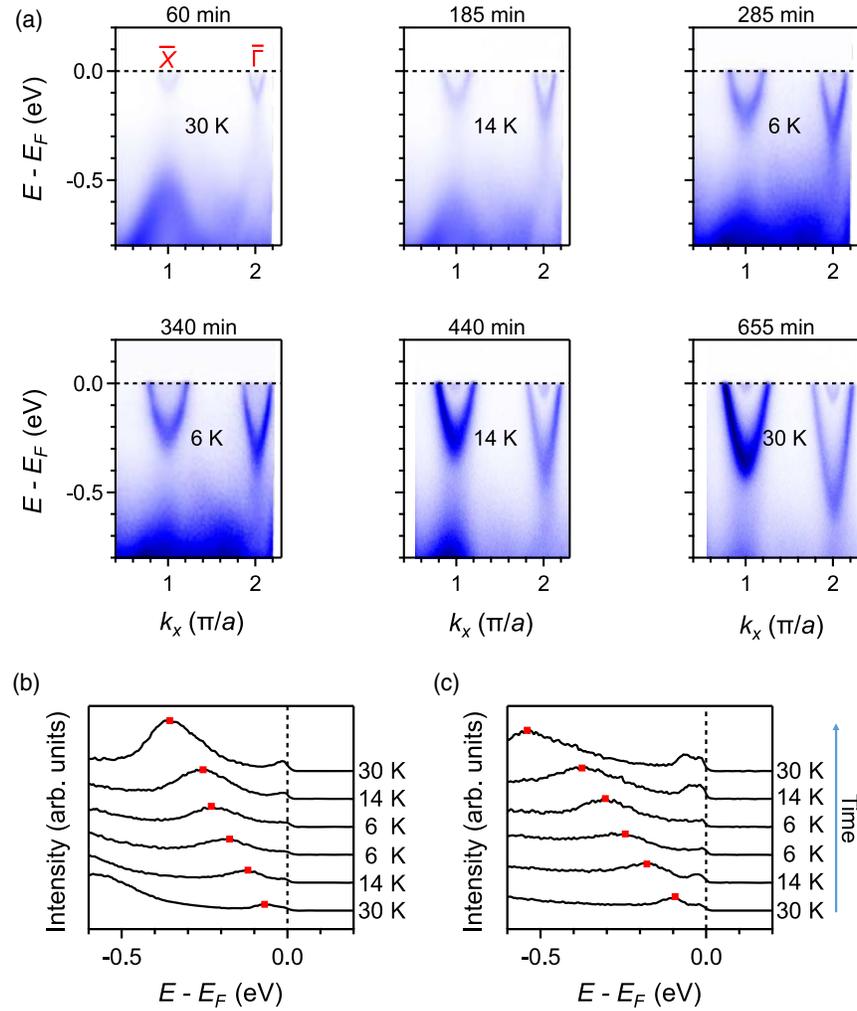


FIG. 6. Time dependence of band dispersions measured on Eu-terminated surfaces under temperature cycling. (a) Intensity plots of the ARPES data collected under temperature cycling in time sequence. (b),(c) Energy distribution curves at the \bar{X} and $\bar{\Gamma}$ points, respectively, showing that energy shifts of the band bottoms are related to time rather than temperature.

APPENDIX B: SURFACE DOPING EFFECTS ON Eu- AND B-TERMINATED SURFACES

Since the cleavage produces polar (001) surfaces with either Eu or B termination, it is expected that the electronic states on the polar surfaces suffer from surface doping. As B has a stronger electronegativity than Eu, Eu has a positive valence state in EuB_6 . The surface doping would cause the electronic states on Eu-terminated surfaces to be more n type compared with those in the bulk, which could explain the electron pockets observed on Eu-terminated surfaces in Figs. 2(d) and 2(e).

As B has a negative valence state in EuB_6 , the surface doping would cause the electronic states on B-terminated surfaces to be more p type compared with those in the bulk. However, the electronic states observed on B-terminated surfaces appear to be either nearly neutral in Fig. 2(i) or

slightly n type in Fig. 7, which seems to contradict the expectation. One possible scenario is that the electronic states on B-terminated surfaces are coherent with the bulk states since they exhibit an obvious k_z dispersion, which significantly reduces the surface doping effects. Another possibility is that the bulk states are slightly n type because Hall effect measurements show a negative Hall coefficient [46–48], and the surface hole doping on B-terminated surfaces tends to neutralize the intrinsic n -type carriers.

The electronic states on B-terminated surfaces are observed to be nearly neutral in six samples and slightly n type in two samples, indicating small fluctuations in the stoichiometry of the samples. In addition, the data in Fig. 7 confirm the semiconducting electronic structure with a narrow band gap in the PM state, which is consistent with the calculations in Fig. 1(d).

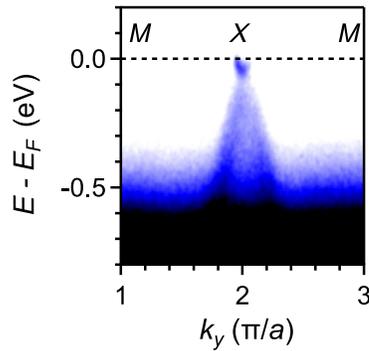


FIG. 7. Intensity plot of the ARPES data along M - X in the k_x - k_y plane with $k_z = 8\pi/a$ at 22 K, showing slightly n -type electronic states on B-terminated surfaces of some samples.

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